

THERMODYNAMIC AND KINETIC CHARACTERIZATION OF VITREOUS EUTECTIC  
GeSe<sub>2</sub>-Sb<sub>2</sub>Te<sub>3</sub> ALLOY\*

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ABSTRACT

For the vitreous alloy of quasi-binary GeSe<sub>2</sub>-Sb<sub>2</sub>Te<sub>3</sub> eutectic composition heat capacities have been measured using differential scanning calorimetry. Enthalpy entropy and Gibbs free energy differences of vitreous and crystalline forms have been deduced and combined with the crystallization kinetic parameters in order to obtain a complete description of this glass.

INTRODUCTION

The phase diagram and glass forming ability of the GeSe<sub>2</sub>-Sb<sub>2</sub>Te<sub>3</sub> system have been reported earlier (ref. 1,2). The eutectic composition and temperature are, respectively, 24 mol% Sb<sub>2</sub>Te<sub>3</sub> and 758±3K. The glass forming region extends from 5 to 30 mol% Sb<sub>2</sub>Te<sub>3</sub>.

In this paper we report on the determination of heat capacity around the glass transition of a chalcogenide glass of nominal composition (GeSe<sub>2</sub>)<sub>76</sub>(Sb<sub>2</sub>Te<sub>3</sub>)<sub>24</sub> and the investigation of the residual values of thermodynamic quantities with respect to the stable state. The results are correlated to kinetic and morphological studies.

EXPERIMENTAL

Glassy samples were prepared by room temperature water quenching of the molten alloys. The thermal behaviour was investigated using a Perkin-Elmer DSC-2 coupled with a data acquisition system. We followed the normal procedure of calibrating measured temperature and energy values.

RESULTS

The DSC traces show the glass transition and two well-defined, through overlapped, exothermic crystallization peaks. In fig.1 the heat capacities are plot-

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\* Work sponsored by Universitat de Barcelona (Ajut Investigació 1983).

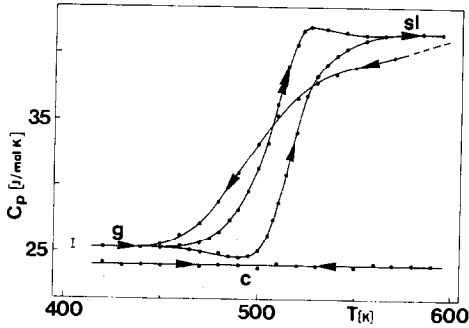


Fig. 1. Heat capacity versus temperature for glassy (g), supercooled liquid (sl) and crystalline (c) phases.

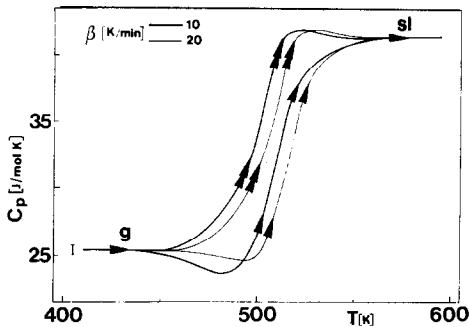


Fig. 2. Heat capacity curves in the glass transition region as a function of scan rate.

ted versus temperature for a glass heated ( $\rightarrow$ ), cooled ( $\leftarrow$ ) and then reheated ( $\leftrightarrow$ ) through the glass transition region and for the crystalline phases, at a rate of  $20 \text{ K min}^{-1}$ . The points represent some of the data obtained by numerical treatment of a typical run and the bar the error for runs on different samples. As seen here a glass obtained by slow cooling has a smaller glass transition temperature  $T_g$  than the as-quenched one. Allowing for the kinetic nature of glass transition there is also an increase of  $T_g$  with heating rate. This fact is illustrated in fig.2 where are presented the results obtained for both as-quenched and slowly cooled glasses, at different scanning rates. Below  $T_g$  the  $C_p$  values agree with the classical limit of  $3R$ . The change of  $C_p$  at the glass transition,  $\Delta C_p$ , is  $16.3 \text{ J K}^{-1} \text{ mol}^{-1}$ , a typical value in many chalcogenide systems (see table 1).

From heat capacity determination enthalpy and entropy values for

glass, supercooled liquid and crystalline phases have been calculated relative to the standard state (crystalline phase at  $298 \text{ K}$ ). Figure 3 shows the entropy changes of as-prepared glasses during heating (path:  $A \rightarrow B \rightarrow C$ ), cooling ( $C \rightarrow D \rightarrow E$ ), and reheating ( $E \rightarrow D \rightarrow C$ ) at a scan rate of  $10 \text{ K min}^{-1}$ . The temperature  $T_0$  at which the entropy difference between the liquid and the crystalline phase vanishes defines, according to Kauzmann (ref. 7), the lower limit for the temperature of glass transition. Furthermore  $T_g$  would approach  $T_0$  in the limit of low cooling or heating rates. We obtain  $T_0 = 457 \text{ K}$ ,  $T_g$  (as-quenched glass) =  $T_1 = 522 \text{ K}$  while  $T_g$  (after cooling and reheating at  $10 \text{ K min}^{-1}$ ) =  $T_2 = 492 \text{ K}$ . That is, extrapolating the liquid state through the glass transition region, the calculated  $T_0$  reproduces the trends of the experimentally determined  $T_g$  values ( $T_1$  and  $T_2$ ). The residual entropy (at  $298 \text{ K}$ ) of the as-quenched glass with respect to the crystalline sample is  $2.0 \text{ J mol}^{-1} \text{ K}^{-1}$ . This value is much smaller than those obtained for

TABLE 1.

Heat capacity change at the glass transition and residual entropy of some glasses

Glass	As <sub>2</sub> S <sub>3</sub>	As <sub>2</sub> Se <sub>3</sub>	As <sub>2</sub> Te <sub>3</sub>	Se	P <sub>4</sub> Se <sub>4</sub>	Ge <sub>25</sub> Sb <sub>15</sub> Se <sub>60</sub>	our Sample
$\Delta C_p$ (J g·at <sup>-1</sup> K <sup>-1</sup> )	13.0 <sup>a</sup>	14.6 <sup>a</sup>	-	-	-	15.9 <sup>b</sup>	16.3
$\Delta S^{\text{res}}$ (J g·at <sup>-1</sup> K <sup>-1</sup> )	6.9 <sup>c</sup>	5.8 <sup>c</sup>	5.8 <sup>d</sup>	7.2 <sup>d</sup>	2.8 <sup>c</sup>	-	2.0
	7.0 <sup>d</sup>	8.5 <sup>d</sup>					

a: from ref. 3; b: from ref. 4; c: from ref. 5; d: from ref. 6

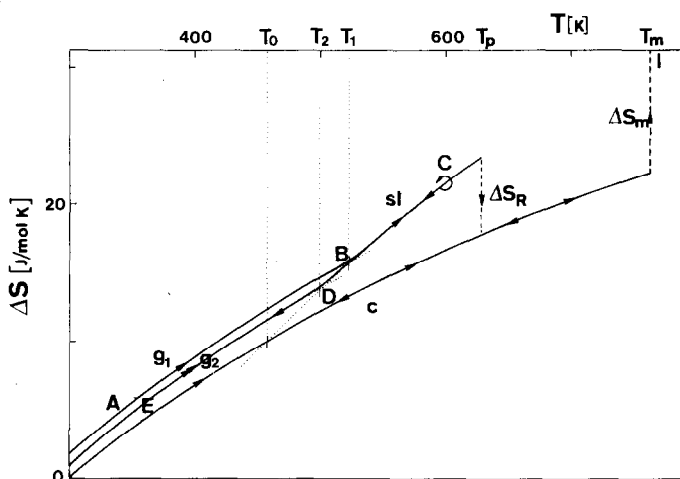


Fig. 3. Plot of entropy versus temperature for glasses ( $g_1, g_2$ ), supercooled liquid (sl), crystalline (c) and liquid (l) phases.  $\Delta S_R$  ( $\Delta S_m$ ) is the crystallization (melting) entropy at temperature  $T_p$  ( $T_m$ ).

other chalcogenide glasses (see table 1). Furthermore, the residual entropy of the slowly cooled glass is about one half of that of the as-quenched one. These results further suggest that both glasses deviate little from the state of an ideal glass.

The heat capacity measurements allow also for the determination of the Gibbs free energy difference,  $\Delta G$ , between supercooled liquid and crystalline solid. Curve 1 of fig. 4 shows the value of  $\Delta G$  obtained from experimental results. Curve 3 gives the plot obtained using Turnbull (ref. 8) formula  $\Delta G = \Delta S_m(T_m - T)$ . A best fit is obtained with Hoffmann (ref. 9) equation (curve 2) which assumes that  $\Delta G$  is  $T/T_m$  times Turnbull's value. This result supports the existence of associations in the supercooled liquid. The Gibbs free energy difference between the supercooled liquid and the crystalline phases at the onset of crystallization is  $\Delta G_c = 1.2 \text{ kJ mol}^{-1}$ . From simple nucleation theory (ref. 10) the activation energy for nucleation is then expected to be  $E_N \sim (T_m \Delta S_m)^3 / (\Delta G_c)^2 \sim 218 \text{ kJ mol}^{-1}$ . The crystallization process occurs in two well-defined, through overlap-

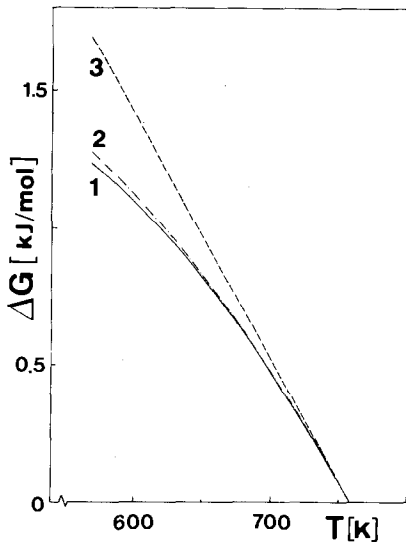


Fig. 4. Measured and estimated temperature dependence of the Gibbs free energy difference between liquid and crystalline phases.

ped peaks. The peak method (ref. 11) has been used to obtain the apparent activation energies. The values deduced are  $150 \pm 5$  and  $230 \pm 7$   $\text{kJ mol}^{-1}$ , respectively, for the first and second peak. Although the preceding calculation of  $E_N$  is an approximate one it accounts for the values obtained for the apparent activation energies. Optical microscopic examinations show that nuclei grow in spherulitic form in the bulk. A surface nucleation density of  $3 \times 10^8 \text{m}^{-2}$  was observed for as-quenched glasses heated just to the onset of crystallization.

#### CONCLUSIONS

Cycled heat capacity measurements around the glass transition demonstrate a "hysteresis-like" behaviour of the eutectic  $\text{GeSe}_2\text{-Sb}_2\text{Te}_3$  vitreous alloy. The glass transition temperature increases with heating rate and decreases for slowly cooled glasses, but neither the value of  $C_p$  for the glass nor its change  $\Delta C_p$  at the glass transition vary. The decrease of residual entropy further evidences the thermal relaxation of the glass towards a more stable state. The Gibbs free energy difference between supercooled liquid and crystalline phases at the onset of crystallization accounts for the effective activation energy of crystallization obtained.

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